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*Paper: Transformability*

## Cosmological constant in gravitational lensing

Consider the Schwarzschild de Sitter Metric,

$$ds^2 = \left(1 - \frac{2GM}{rc^2} - \frac{\Lambda r^2}{3}\right) c^2 dt^2 - \left(1 - \frac{2GM}{rc^2} - \frac{\Lambda r^2}{3}\right)^{-1} dr^2 - r^2(d\theta^2 + \sin^2 \theta d\phi^2). \quad (1)$$

The constant term  $\frac{2GM}{c^2}$  is recognized as the Schwarzschild radius ( $r_s$ ), and typically it is replaced by a constant term  $2m$ , where  $m = \frac{1}{2}r_s = \frac{GM}{c^2}$  and then the equation (1) can be written as follows.

$$ds^2 = \left(1 - \frac{2m}{r} - \frac{\Lambda r^2}{3}\right) c^2 dt^2 - \left(1 - \frac{2m}{r} - \frac{\Lambda r^2}{3}\right)^{-1} dr^2 - r^2(d\theta^2 + \sin^2 \theta d\phi^2). \quad (2)$$

$\Lambda$  is the cosmological constant.

The null-geodesic equation in Schwarzschild-de Sitter metric can be written as,

$$\frac{E^2}{c^2} - l^2 u'^2 - l^2 u^2 + 2ml^2 u^3 + \frac{\Lambda l^2}{3} = 0, \quad (3)$$

where  $E$  is the energy,  $l$  is the orbital angular momentum,  $\Lambda$  is the cosmological constant,  $u = \frac{1}{r}$  and

$$u' = \frac{du}{d\phi}.$$

Differentiating (3) with respect to  $\phi$ ,

$$u'(u'' + u - 3mu^2) = 0. \quad (4)$$

Neglecting the solution,  $u' = 0$  which implies  $u = \text{constant}$ , the equation of a light ray trajectory can be written as,

$$u'' + u = 3mu^2. \quad (5)$$

The zeroth order solution and the first order solution of the equation (5) that represent the light ray trajectory are respectively given below.

$$u_0 = \frac{1}{r_0} \cos \phi \quad (6)$$

$$u = \frac{1}{r_0} \cos \phi - \frac{\varepsilon}{3r_0^2} \cos^2 \phi + \frac{2\varepsilon}{3r_0^2} \quad (7)$$

where  $\varepsilon = 3m$ .

In general, in the literature, it is assumed that (7) is a solution of equation (3) without considering the limitations imposed. In this paper we discuss conditions under which (7) is a solution of equation (3).

Now the orbital angular momentum,  $l = pr_0$  where  $p$  is the linear momentum.

The linear momentum,  $p = \frac{E}{c}$ .

Therefore,

$$l = \frac{E}{c} r_0. \tag{8}$$

Substituting (7) and (8) in (3), we have,

$$\begin{aligned} & \frac{E^2}{c^2} - l^2 \left[ -\frac{1}{r_0} \sin \phi + \frac{2\varepsilon}{3r_0^2} \sin \phi \cos \phi \right]^2 - l^2 \left[ \frac{1}{r_0} \cos \phi - \frac{\varepsilon}{3r_0^2} \cos^2 \phi + \frac{2\varepsilon}{3r_0^2} \right]^2 \\ & + \frac{2\varepsilon}{3} l^2 \left[ \frac{1}{r_0} \cos \phi - \frac{\varepsilon}{3r_0^2} \cos^2 \phi + \frac{2\varepsilon}{3r_0^2} \right]^3 + \frac{\Lambda l^2}{3} = 0. \end{aligned} \tag{9}$$

By simplifying the above equation and since  $l \neq 0$  we obtain the following equation,

$$\begin{aligned} & 2m \left[ \frac{8\varepsilon^3}{27r_0^6} - \frac{4\varepsilon^3}{9r_0^6} \cos^2 \phi + \frac{2\varepsilon^3}{9r_0^6} \cos^4 \phi - \frac{\varepsilon^3}{27r_0^6} \cos^6 \phi + \frac{\varepsilon^2}{3r_0^5} \cos^5 \phi - \frac{4\varepsilon^2}{3r_0^5} \cos^3 \phi + \frac{4\varepsilon^2}{3r_0^5} \cos \phi \right] + \frac{\Lambda}{3} = 0 \\ & \Lambda = -18m^2 \left[ \frac{8m^2}{3r_0^6} - \frac{4m^2}{r_0^6} \cos^2 \phi + \frac{2m^2}{r_0^6} \cos^4 \phi - \frac{m^2}{3r_0^6} \cos^6 \phi + \frac{m}{r_0^5} \cos^5 \phi \right. \\ & \left. - \frac{4m}{r_0^5} \cos^3 \phi + \frac{4m}{r_0^5} \cos \phi - \frac{2}{3r_0^4} + \frac{2}{r_0^4} \cos^2 \phi - \frac{1}{2r_0^4} \cos^4 \phi \right]. \end{aligned} \tag{10}$$

From (10) it is clear that the solution given by (7) of equation (3) is valid only if  $\Lambda$  is a constant of order  $m^2$ , and as we neglect terms of order 2 and above we are justified in assuming (7) as a solution of equation (3). However, it turns out that this particular solution is valid only if  $\Lambda$  is a constant of order 2 or more in  $m$ . If  $\Lambda$  is a non zero constant and of order one in  $m$ , the solution (7) is not valid and we have to seek other solutions.

**References**

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